Electric flux: is defined as total number of electric field lines passing from any area, denoted by mathematically it is defined as integral of **E** and  $dS \implies \phi = \int E \cdot dS$  with units of Newton meter<sup>2</sup>/Coulomb **Remember area vector S** is perpendicular to the surface of area under consideration.

Gauss Law: states that surface integral of electric field produced by any source over any closed surface sing a volume V is  $1/\varepsilon_0$  times total charge contained inside surface. **S** enclosing a volume V is  $1/\varepsilon_0$  times total charge contained inside surface.

It is to note that electric charges outside Gaussian surface do not contribute to electric field.

Derivation of Coulomb law, E due to linear charged & E due to a charged sphere (Hollow/Conductor/uniformly Charged)

Basics: in all the derivations we get a Gaussian surface symmetrical to charge. Find out direction of **E** and S at various points on surface. Whenever **E** is perpendicular to surface **E**.dS is zero and if **E** is  $\|$  to  $\|$  **E.** S becomes **EdS.** After calculating flux we find the charge in given Gaussian surface. And use  $\int$ **E**.d**S** =  $Q/\epsilon_0$ 



In  $1<sup>st</sup>$  case we take point charge Q placed at center, then taking sphere as shown in Figure 1. Considering point 'A' near the surface we see that  $\mathbf{E}'$  and 'd**S'** are in same direction thus  $t$ ussian surface (radius r) we draw as

- $=\oint E.dS = E \oint dS$  (E is constant)  $\phi = \oint E \cdot dS = \oint E \cdot dS \cos \theta$
- $=$  E.4 $\pi r^2$

Now as per Gauss law it should be equal  $\circ$  $e_0 \rightarrow L - 4\pi \varepsilon_0 r^2$ 2 r Q 4  $E.4\pi r^2 = \frac{Q}{r} \Rightarrow E = \frac{1}{r}$  $\phi = E.4\pi r^2 = \frac{Q}{\epsilon_0} \Rightarrow E = \frac{1}{4\pi\epsilon}$ Now as per Gauss law it should be equal to contained in volume i.e. Q. equating both we g<br>  $\phi = E.4\pi r^2 = \frac{Q}{\epsilon_0} \Rightarrow E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$ <br>
If we place charge Q' at point A the force is given by **F** = Q'**E** placing the e contained in volume i.e. Q. equating both we get

QQ' 4  $F = \frac{1}{4\pi\epsilon}$ 

which is Coulomb's law. In integration we have taken distance from center, thus  $s$  and we have taken  $\sum_{n=1}^{\infty}$  or  $\sum_{n=1}^{\infty}$  of the integral with assumption that due to symmetry each point is same



In Second case for a linear conductor  $(\lambda$  Coulomb/meter) charge we take Gaussian surface as a cylinder (radius r & length L) with axis collinear with the charge.

Drawing different vectors at different places for **E** and **S** we see that at the curved surface **E** and **S** are parallel to each other, while at plain surfaces they are perpendicular to each other.

Thus only **E** & **S** which are on curved surface contribute to the calculation of flux. Mathematically

$$
\phi = \oint EdS = \int_{CSA} EdS + \int_{Eads} EdS
$$
  
\n
$$
= \int_{CSA} EdS - \int_{CSA} EdS - \int_{Eads} EdS
$$
  
\n
$$
= E2\pi L
$$
  
\n
$$
= E2\pi L
$$
  
\n
$$
E2\pi L = \frac{2L}{\epsilon_0} \Rightarrow E = \frac{\lambda}{2\pi\epsilon_0}
$$
  
\n
$$
E2\pi L = \frac{\lambda L}{\epsilon_0} \Rightarrow E = \frac{\lambda}{2\pi\epsilon_0}
$$
  
\n
$$
\oint
$$
  
\nIn this case we try to get *E* due to hollow or charge *Q* and *Q* and *Q* is a factor (radial to charge) in both the cases charge is residing at the periphery. So we take 3 different gaussian super  
\n(ii)  $r = R$  sphere 'A'  
\n(iii)  $r > R$  sphere 'B'  
\n(iii)  $r > R$  sphere 'B'  
\n
$$
\oint = \oint EdS = \frac{Q}{\epsilon_0}
$$
 but since c)  
\n $\oint = \oint EdS = \frac{Q}{\epsilon_0} = 0$  as d $\Rightarrow \epsilon$  *Q*  
\nCase (ii)  $\phi = \oint EdS = \frac{Q}{\epsilon_0}$  but since c)  
\n $\theta = E4\pi R^2$   
\nCase (iii)  $\oint = \frac{Q}{\epsilon_0} = 0$  as d $\Rightarrow \epsilon$  *Q*  
\nCase (iii)  $\oint = d\pi c_0$   $\frac{1}{R^2}$   
\nCase (iv)  $\Phi = \oint EdS = Ef dS$  also  
\n
$$
\oint = \frac{1}{4\pi c_0} \frac{Q}{R^2}
$$
  
\nCase (iii)  $\oint = \frac{1}{4\pi c_0} \frac{Q}{R^2}$   
\n
$$
\oint = \frac{1}{4\pi c_0} \frac{Q}{R^2}
$$



Lastly we take derivation of uniformly charged sphere, where charge 'Q' is uniformly distributed over the full sphere.

Obviously it has to be insulator. Thus the derivation for  $r \ge R$  is same as for conductor (hollow or solid sphere). Charge per unit volume( $\rho$ ) = Q $\sqrt{2\pi R^3}$  $\overline{\mathsf{R}}$ B The only change is in the  $E$  inside the sphere, as in this case some charge is there, while in previous cases there was no charge. Applying Gauss in this case  $\left( \begin{array}{c} \sqrt{2} \end{array} \right)$ Which is equal to charge in Guassian sphere 3 3  $2$  Qr<sup>3</sup>  $\sim$  $\angle$ Thus  $\phi = \frac{Qr}{r}$  $E.4\pi r^2 = \frac{Qr}{r^2}$  $E = \frac{1}{1}$ Qr 3 Thus  $\phi = \frac{Q_1}{\epsilon_0 R^3} \Rightarrow E.4\pi r^2 = \frac{Q_1}{\epsilon_0 R^3} \Rightarrow E = \frac{1}{4\pi \epsilon_0 R^3}$ 3 3  $_0$ R<sup>3</sup>  $_{0}R^{3}$   $\rightarrow$   $L _{4\pi \varepsilon_{0}}$   $\mathbb{R}^{3}$ 4 R R R R 3 Tas `dS  $\ddot{}$ dS 2 dS  $\ddot{}$ dS Considering the case of infinitely large dimensioned charged sheet having charge per unit area as  $\sigma$  Coul.m<sup>2</sup>. We take cylindrical Guassian surface with axis perpendicular to the plane of sheet. It is clear that at curved surface are **E** and d**S** are perpendicular to each other. Thus CSA does not contribute to flux, but only circular sheet at the extreme ends contribute to flux as **E** and d**S** are parallel in this case.

 $\sigma$  $\phi = \oint E \cdot dS = E \cdot 2A =$ whet interesting to note that in present case  $\bf{E}$  is independent of distance  $rac{371}{\epsilon_0} \Rightarrow E = \sum$ e from sheet. *Why*?  $\ddot{}$  $\bf{B}$  $C_A$  $\overline{C}$  $\mathbf{A}$ 

## E due to Charged Spherical SHELL/CONDUCTING SPHERE

